





AN ACTION APPROACH TO NODAL AND LEAST ENERGY NORMALIZED SOLUTIONS FOR NONLINEAR SCHRÖDINGER EQUATIONS Colette De Coster⁽¹⁾, Simone Dovetta⁽²⁾, <u>Damien Galant^(1, 3), Enrico Serra⁽²⁾</u>

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The nonlinear Schrödinger equation

We are interested in nonzero solutions of the PDE

$$\int -\Delta u + \lambda u = |u|^{p-2} u \quad \text{in } \Omega,$$



 $\begin{cases} u = 0 & \text{in } \partial\Omega, \\ u = 0 & \text{in } \partial\Omega, \end{cases}$ where p > 2 and λ are real parameters and $u : \Omega \to \mathbb{R}$ and $\Omega \subseteq \mathbb{R}^N$ is an open domain. For a given λ , solutions of (NLS) correspond to critical points of the *action functional* $J_{\lambda} : H_0^1(\Omega) \to \mathbb{R}$, defined by (NLS)

 $J_{\lambda}(u) := \frac{1}{2} \int_{\Omega} \|\nabla u\|^2 \,\mathrm{d}x + \frac{\lambda}{2} \int_{\Omega} |u|^2 \,\mathrm{d}x - \frac{1}{p} \int_{\Omega} |u|^p \,\mathrm{d}x.$

Action ground states

The functional J_{λ} is not bounded from below on $H_0^1(\Omega)$. Indeed, if $u \neq 0$, then

$$J_\lambda(tu)=rac{t^2}{2}\|u'\|_{L^2(\Omega)}^2+rac{\lambda t^2}{2}\|u\|_{L^2(\Omega)}^2-rac{t^p}{p}\|u\|_{L^p(\Omega)}^p \xrightarrow[t o\infty]{} -\infty.$$

A typical way to recover boundedness consists in introducing the Nehari manifold \mathcal{N}_{λ} :

$$\begin{split} \mathcal{N}_{\lambda} &:= \Big\{ u \in H^1(\Omega) \setminus \{0\} \mid J_{\lambda}'(u)u = 0 \Big\} \\ &= \Big\{ u \in H^1(\Omega) \setminus \{0\} \mid \|\nabla u\|_{L^2(\Omega)}^2 + \lambda \|u\|_{L^2(\Omega)}^2 = \|u\|_{L^p(\Omega)}^p \Big\}. \end{split}$$

One then defines the *action ground state level*

$$\mathcal{J}(\lambda) := \inf_{u \in \mathcal{N}_{\lambda}} J_{\lambda}(u)$$

It is standard to show (see e.g. [8]) that, when Ω is bounded, $p < 2^*$ and $\lambda > -\lambda_1(\Omega)$, minimizers of the above problem exist and are *least action solutions of the problem*. They have a constant sign since one may replace u by |u| in the minimization problem. $M_p^{nod}(\Omega) := \left\{ \|u\|_{L^2(\Omega)}^2 \mid u \in \mathcal{N}_{\lambda}^{nod} \text{ and } J_{\lambda}(u) = \mathcal{J}^{nod}(\lambda) \text{ for some } \lambda \in \mathbb{R} \right\}$ be the set of masses of all action ground states and nodal action ground states. Then (i) if $p < 2 + \frac{4}{N}$, then $M_p(\Omega) = M_p^{nod}(\Omega) = (0, \infty);$ (ii) if $p = 2 + \frac{4}{N}$, then there exist $0 < \mu_p, \mu_p^{nod} < \infty$ such that $(0, \mu_p) \subseteq M_p(\Omega) \subseteq (0, \mu_p] \text{ and } (0, \mu_p^{nod}) \subseteq M_p^{nod}(\Omega) \subseteq (0, \mu_p^{nod}];$ (iii) if $p > 2 + \frac{4}{N}$, then there exist $0 < \mu_p, \mu_p^{nod} < \infty$ such that $M_p(\Omega) = (0, \mu_p] \text{ and } M_p^{nod}(\Omega) = (0, \mu_p^{nod}].$

Theorem. Let $\Omega \subset \mathbb{R}^N$ be open and bounded, and either (i) $p < 2 + \frac{4}{N}$ and $\mu > 0$; or (ii) $p = 2 + \frac{4}{N}$ and $\mu < 2\mu_N$, where $\mu_N := 2 \inf_{u \in \mathcal{N}_1(\mathbb{R}^N)} J_1(u)$; or (iii) $p > 2 + \frac{4}{N}$, Ω is star-shaped, and μ is small enough. Then there exists a *least energy normalized nodal solution* with mass μ .

An idea of the techniques

We make a strong use of the *convex duality* between the action and the energy levels discovered in [4].

Nodal action ground states

All sign-changing solutions of the problem belong to the nodal Nehari set

 $\mathcal{N}_{\lambda}^{nod} := \left\{ u \in H_0^1(\Omega) \mid u^{\pm} \in \mathcal{N}_{\lambda} \right\}$

where $u^+ = \max(u, 0)$ and $u^- = \min(u, 0)$ are the positive and negative parts of u. Then, one considers the level

 $\mathcal{J}^{nod}(\lambda) := \inf_{u \in \mathcal{N}_{\lambda}^{nod}} J_{\lambda}(u).$

When Ω is bounded, $p < 2^*$ and $\lambda > -\lambda_2(\Omega)$, minimizers of the above problem exist (see [1]) and are *least action nodal solutions of the problem*.

Normalized solutions

A normalized solution to (NLS) is a solution whose L^2 -norm (usually called the mass) is prescribed a priori, whereas λ is an unknown of the problem. They correspond to constrained critical points of the energy functional $E: H_0^1(\Omega) \to \mathbb{R}$

$$E(u) := \frac{1}{2} \|\nabla u\|_{L^2(\Omega)}^2 - \frac{1}{p} \|u\|_{L^p(\Omega)}^p$$

on the L^2 -sphere

$$\mathcal{M}_{\mu} := \left\{ u \in H_0^1(\Omega) \mid \|u\|_{L^2(\Omega)}^2 = \mu \right\}$$

In particular, we show that if $\lambda_*, \lambda_*^{nod} \in \mathbb{R}$ are local minimizers of the maps

 $\lambda \mapsto \mathcal{J}(\lambda) - \frac{\mu}{2}\lambda \quad \text{or} \quad \lambda \mapsto \mathcal{J}^{nod}(\lambda) - \frac{\mu}{2}\lambda,$

then action ground states in \mathcal{N}_{λ_*} and nodal action ground states in $\mathcal{N}_{\lambda_*^{nod}}^{nod}$ have mass μ . We then show that such minimizers exist under suitable assumptions on the masses.

A counterintuitive result

We take $p = 2 + \frac{4}{N}$ and we consider the ball. Noris, Tavares and Verzini have shown in [6] that the set $\left\{ \|u\|_{L^2(\Omega)}^2 \mid u \text{ positive solution of (NLS) for some } \lambda \in \mathbb{R} \right\}$

is equal to $(0, \mu_N)$. Our results imply that, for $\mu \in [\mu_N, 2\mu_N)$, there exist least energy normalized nodal solutions with mass μ . Thus...

Least energy solutions may be nodal!

References

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λ arising then as a Lagrange multiplier.

Least energy normalized solutions, i.e. functions $u \in \mathcal{M}_{\mu}$ solving (NLS) and satisfying

$E(u) = \inf \left\{ E(v) \mid v \in \mathcal{M}_{\mu} \text{ solves (NLS) for some } \lambda \in \mathbb{R} \right\}$

are particularly interesting. When $p < 2 + \frac{4}{N}$, least energy solutions can be found (see e.g. [2]) by solving the minimization problem $\inf_{u \in \mathcal{M}_{\mu}} E(u)$.

When $p > 2 + \frac{4}{N}$, the energy E is unbounded from below on \mathcal{M}_{μ} for every μ , and different approaches are needed (e.g. of mountain-pass type, see [5]). This problem received a lot of interest on \mathbb{R}^N but, so far, much less on bounded domains (see e.g. [6] and [7] however).

Two questions

How to find least energy normalized solutions in the L^2 -supercritical regime?

How to find normalized nodal solutions?

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